

# Quantum Mechanics

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## Chapter 3

# The Time-Independent Schrödinger Equation

In the Schrödinger Equation

$$\left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x, t) \right] \Psi(x, t) = i\hbar \frac{\partial \Psi(x, t)}{\partial t}$$

the potential often does not depend on time; i.e.  $V(x, t) = V(x)$ . An example of this would be the infinite square well potential.

$$V(x, t) = \begin{cases} 0 & -\frac{a}{2} < x < \frac{a}{2} \\ \infty & |x| \geq \frac{a}{2} \end{cases} \quad (a > 0)$$

When  $V(x, t) = V(x)$ , we can express the multivariable function  $\Psi(x, t)$  as a product of a function of  $x$  and a function of  $t$  as follows.

$$\Psi(x, t) = \psi(x)\phi(t)$$

This technique is called “separation of variables”. Let us substitute the above  $\Psi$  into the Schrödinger Equation.

$$\begin{aligned} \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V(x) \right] \psi(x)\phi(t) &= i\hbar \frac{\partial \psi(x)\phi(t)}{\partial t} \\ \implies -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x)\phi(t) + V(x)\psi(x)\phi(t) &= i\hbar \frac{\partial \psi(x)\phi(t)}{\partial t} \\ \implies -\frac{\hbar^2}{2m} \phi(t) \frac{\partial^2}{\partial x^2} \psi(x) + V(x)\psi(x)\phi(t) &= \psi(x) i\hbar \frac{\partial \phi(t)}{\partial t} \end{aligned}$$

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$$\implies \phi(t) \left[ -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \psi(x) + V(x)\psi(x) \right] = \psi(x) i\hbar \frac{\partial \phi(t)}{\partial t}$$

Dividing through by  $\psi(x)\phi(t)$ , we get

$$\underbrace{\frac{1}{\psi(x)} \left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) + V(x)\psi(x) \right]}_{H(x)} = \underbrace{i\hbar \frac{1}{\phi(t)} \frac{d\phi(t)}{dt}}_{K(t)};$$

where we changed the partial derivatives to total derivatives as we now have one-variable functions.

Generally speaking,

$$H(x) = K(t) \text{ for all pairs } (x, t)$$

means that both  $H(x)$  and  $K(t)$  are a constant. Let  $H(x) = K(t) = G$  (a constant). Then,

$$K(t) = G \iff i\hbar \frac{1}{\phi(t)} \frac{d\phi(t)}{dt} = G \iff \frac{d\phi(t)}{dt} = \frac{1}{i\hbar} G \phi(t) = \frac{-i}{\hbar} G \phi(t).$$

Therefore,

$$\phi(t) = A e^{-iGt/\hbar}.$$

Compare this with the time-dependent part of the travelling wave

$$e^{i(kx - \omega t)} = e^{ikx} e^{-i\omega t}.$$

We realize that  $\frac{G}{\hbar} = 2\pi \frac{G}{h}$  “should be”  $\omega$ .

$$2\pi \frac{G}{h} = \omega \iff G = \frac{\omega h}{2\pi} = \frac{2\pi \nu h}{2\pi} = h\nu = E.$$

We now have  $G = E$ , and

$$\phi(t) = e^{-iEt/\hbar}.$$

This also implies

$$H(x) = \frac{1}{\psi(x)} \left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) + V(x)\psi(x) \right] = E \implies -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) + V(x)\psi(x) = E\psi(x);$$

where  $E$  is the total energy.

$$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \psi(x) + V(x)\psi(x) = E\psi(x) \text{ or } \left[ -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right] \psi(x) = E\psi(x)$$

is the Time-Independent Schrödinger Equation.

Note here that the full wavefunction is given by

$$\Psi(x, t) = \psi(x)e^{-iEt/\hbar}.$$

In the rest of this course, we will solve the Time-Independent Schrödinger Equation for various systems, culminating in a solution for the hydrogen atom.